Ph 20.6 – Numerical Computations with Mathematica: “The Grapefruit Problem”

Due: Week 9

This Assignment

The previous assignment was meant to familiarize you with the capabilities of Mathematica and get you used to its syntax and method of programming. This assignment combines some of the symbolic capabilities of Mathematica which you have already seen with its numerical tools. So far, the Ph20 assignments have dealt with particles undergoing free motion in a potential. This assignment has as its subject the motion of particles when a dissipative force (drag) is present. Its ultimate goal is to find the velocity-minimizing trajectory to launch a projectile a given distance.

You should be aware that this assignment is probably the most challenging one in Ph 20. Below we sketch a series of steps for approaching the problem: first solving for the drag-free situation given initial conditions, then adding the influence of drag, and finally constructing a hierarchy of functions to find the initial conditions which lead to the optimal trajectory. The hints at the end of the writeup, your Mathematica experience from the previous assignment, and built-in function documentation should give you everything you need to carry out these steps. Your responsibility, beyond these basic steps, is to ensure you are correctly capturing the influence of drag, to make informative and illustrative plots, and to come up with an efficient method for iterating to find the optimal trajectory. Be sure to sanity-check your solutions at each stage using plots and physical intuition, as we have emphasized in the first three assignments of this course.

Where did that grapefruit come from?

Rumor (history) has it that Caltech undergrads used a kerosene cannon to lob grapefruits onto Pasadena City College, at a distance of $\sim 1000$ m. Of course, Caltech students would never do something like this, at least not the students we get these days, but the physics of this venerable legend is a good problem to feed to Mathematica.

Mathematically (and Mathematically), the problem is that of a projectile subject to gravity, and to the resistance of air (atmospheric drag). Leaving apart drag for the moment, the equations of motion should be quite familiar to you:

$$\frac{dv_x}{dt} = 0, \quad \frac{dv_y}{dt} = -g;$$

$$v_x = \frac{dx}{dt}, \quad v_y = \frac{dy}{dt};$$

which (as you know well) can be solved to yield

$$x = x_0 + v_{x0}t, \quad y = y_0 + v_{y0}t - \frac{1}{2}gt^2.$$  \hspace{1cm} (3)

Of course, real projectiles don’t move quite so simply, because of drag. It turns out that the appropriate expression for the drag force in the case of spherical projectiles, with reasonable velocities, is (see Box 1)

$$F_{drag} \sim -\frac{1}{2} \rho r^2 v^2.$$

1
As you can imagine, the physics of atmospheric drag is very complicated. Enter approximations! We know for a fact that the drag force is null for an object at rest, and that it grows with velocity. Under certain mathematical assumptions (of smoothness, for instance), we can then write

$$F_{\text{drag}} = -B_1 v - B_2 v^2 - B_3 v^3 - \cdots$$

The linear term (and other odd terms) vanish because drag does not depend on the sign of the velocity. For reasonable velocities, it turns out that the quadratic term is dominant. The resulting differential equation for the velocity is

$$\frac{dv}{dt} = \frac{P}{mv} - \frac{B_2 v^2}{m}.$$  

We estimate the coefficient $B_2$ by the following argument: to overcome atmospheric drag, the projectile must push out of the way the volume of air directly in front of it. In a time $\Delta t$, the mass of air moved is $m_{\text{air}} \approx \rho A v \Delta t$, where $\rho$ is the density of air, and $A$ is the projectile’s *frontal area*. This air is given a velocity of order $v$, and therefore a momentum $m_{\text{air}} v$ over a time $\Delta t$. It follows that the instantaneous force exerted by the projectile on the air (and therefore, because of Newton’s third law, by the air on the projectile as a drag force) is approximately

$$F_{\text{drag}} = -\rho A v^2.$$  

For a spherical projectile, one finds empirically that the coefficient $r^2 / 2$ works better than the nominal area $4\pi r^2$.

Box 1: A justification of the quadratic dependence of atmospheric drag.

where $\rho$ is the density of air at sea level, or approximately 1.3 kg/m$^3$, and where $r$ is the radius of the projectile. This drag force can then be incorporated into the equations of motion,

$$\frac{dv_x}{dt} = -\frac{|F_{\text{drag}}| v_x}{m}, \quad \frac{dv_y}{dt} = -g - \frac{|F_{\text{drag}}| v_y}{m}, \quad \text{where } v = \sqrt{v_x^2 + v_y^2}. \quad (5)$$

These equations are much harder to solve by pencil and paper; we will use Mathematica.

The Assignment

1. Using Mathematica, prove the well known result that, in the absence of drag, the optimal firing angle (the angle that yields the longer range for a given initial velocity) is $45^\circ$. Then compute the velocity components necessary to reach PCC from Caltech using the optimal firing angle, still not including drag.

2. In Mathematica, write a routine to integrate numerically the equations for motion without drag, and verify the above result. The function to use for this is `NDSolve[]`; see the hints below if Mathematica’s documentation isn’t sufficiently clear. Superimpose several plots of the trajectory, with the same initial velocity but different firing angles, to show visually that $45^\circ$ is optimal. This is another Ph20 Beautiful Plot™: take full advantage of the optional parameters you can use in `Plot[]` or `ParametricPlot[]` to arrange it so that it makes your point as clearly and boldly as possible.

3. In Mathematica, assume the grapefruit has mass 0.5 kg and radius 0.05 m. (If you’re wondering why the mass doesn’t cancel out of the equations, go back and read the previous section again!) Include the acceleration due to drag in the equations of motion, and now predict where the grapefruit will land if you use the initial velocity you have just found.
4. In Mathematica, try increasing the initial velocity components and changing the firing angle until you can reach the range calculated in the drag-free case. Find the optimal firing angle (an approximate solution obtained by trial and error is acceptable; we’ll do better in the next step). Show a visual comparison of trajectories with the same range (Caltech to PCC), but different firing angles and correspondingly different initial velocities. Again, make this graph beautiful and information-dense.

5. To go beyond trial and error, implement the following hierarchy of Mathematica functions (if needed, see the hints below and/or enlist the substantial help of your TA):

(a) write a function that returns the time when the grapefruit lands \((y = 0)\) as a function of initial angle and speed;

(b) using the result of item a, write another function that returns the range \((x \text{ at } y = 0 \text{ and } t \neq 0)\), given the initial angle and speed, and the corresponding landing time;

(c) using the results of items a and b, write a function that iterates to find the initial velocity required to obtain a given range, for a certain initial angle;

(d) using the results of items a to c, write a function that iterates to find the angle with the least initial velocity required to yield a certain range.

Mathematica Hints for Part 2

How do we solve differential equations in Mathematica? First we define a list containing the equations. Mind the different “equal” signs: “=” is an assignment operator (...let eqs be the equation...), whereas “==” is the test operator used to state that the left side of the equation equals its right side.

\[
eqs = \{x''[t]==0, \ y''[t]==-g\}
\]

Notice the “apostrophe” syntax used to denote derivatives. Let’s then define the initial conditions, say

\[
ini = \{x[0]==0, \ y[0]==0, \ x'[0]==10, \ y'[0]==10\}
\]

The system of differential equations can be solved with the NDSolve Mathematica function, which has the following syntax:

\[
\text{rules} = \text{NDSolve}[\text{Join}[\text{eqs},\text{ini}],\{x,y\},\{t,\text{tinit, tend}\}]
\]

When you invoke NDSolve, the parameters of your equation (in this case, \(g\)) and the initial time \(\text{tinit}\) and final time \(\text{tend}\) should have numerical values. You can assign these beforehand, or enclose NDSolve in a function and pass the parameters as arguments. The result returned by NDSolve is not a function, but rather a Mathematica rule, which you should be familiar with from the last assignment. You can however define functions from the rule:

\[
xx[t_] := x[t] /. \text{rules[[1]]};
\]

\[
yy[t_] := y[t] /. \text{rules[[1]]}
\]

The “[[1]]” is needed because NDSolve returns the rules enclosed in a list (the braces), with the first element denoted by 1 instead of 0 like in Python, and we have to get rid of that first. You can then plot \(xx[t]\) and \(yy[t]\) normally, with Plot, and evaluate them directly, as in “\(xx[2]\)".
Mathematica Hints for Part 5

If you are using NDSolve as suggested above, you will need to re-solve the differential equation each time you specify a new set of initial conditions ($\theta$ or $v$). Starting with part (c) you will need to loop over different initial conditions to find the required velocity or angle. This sort of loop is straightforward in Python but unfortunately less so in Mathematica, where everything is a function. The best place to start is with Module[], which lets you set up variables which are used only within a function. You can then do a While[] or For[] loop inside the module as appropriate.

The time it takes for the function in (d) to run can vary substantially, from a few seconds to hours, depending on how you write it. Some tips:

- Remember that many functions, including NDSolve[] and FindRoot[], return replacement rules rather than direct numerical output. You will need to extract the number from the output to compare it to something else in a loop condition.

- NDSolve[] is the most computationally expensive function, so try to call it as few times per loop as necessary.

- You may be tempted to start with $v = 0$ and increment by 1 each time until you reach the required range. This will result in very slow functions—imagine the value of $v$ required to get very far when $\theta = 0.01$, say. Consider starting with large angle and velocity increments and having your function automatically switch to a smaller increment once the approximate value has been found. Or look up and implement a binary search.